

QUANTUM SPACE-TIME

FOAM

and topological

STRINGS

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QUANTUM THEORY of GRAVITY

$$Z \sim \sum \int \mathcal{D}g \exp(-S(g))$$

geometries
+
topologies

STRING THEORY

$$Z = \exp \left[\hbar^{-2} \text{circle} + \text{torus} + \hbar^2 \text{pair of pants} + \dots \right]$$

+ non-perturbative effects

$$\hbar = g_s = e^{\phi}$$

IS THERE A WAY TO
SEE fluctuating
GEOMETRY

IN STRING PERTURBATION THEORY?

HARDLY, IF WE STAY AT
FINITE ORDER OF
STRING PERTURBATION
THEORY



gravitons - "small"
deviations
from background

TOPOLOGICAL STRING IS
THE ONLY EXAMPLE

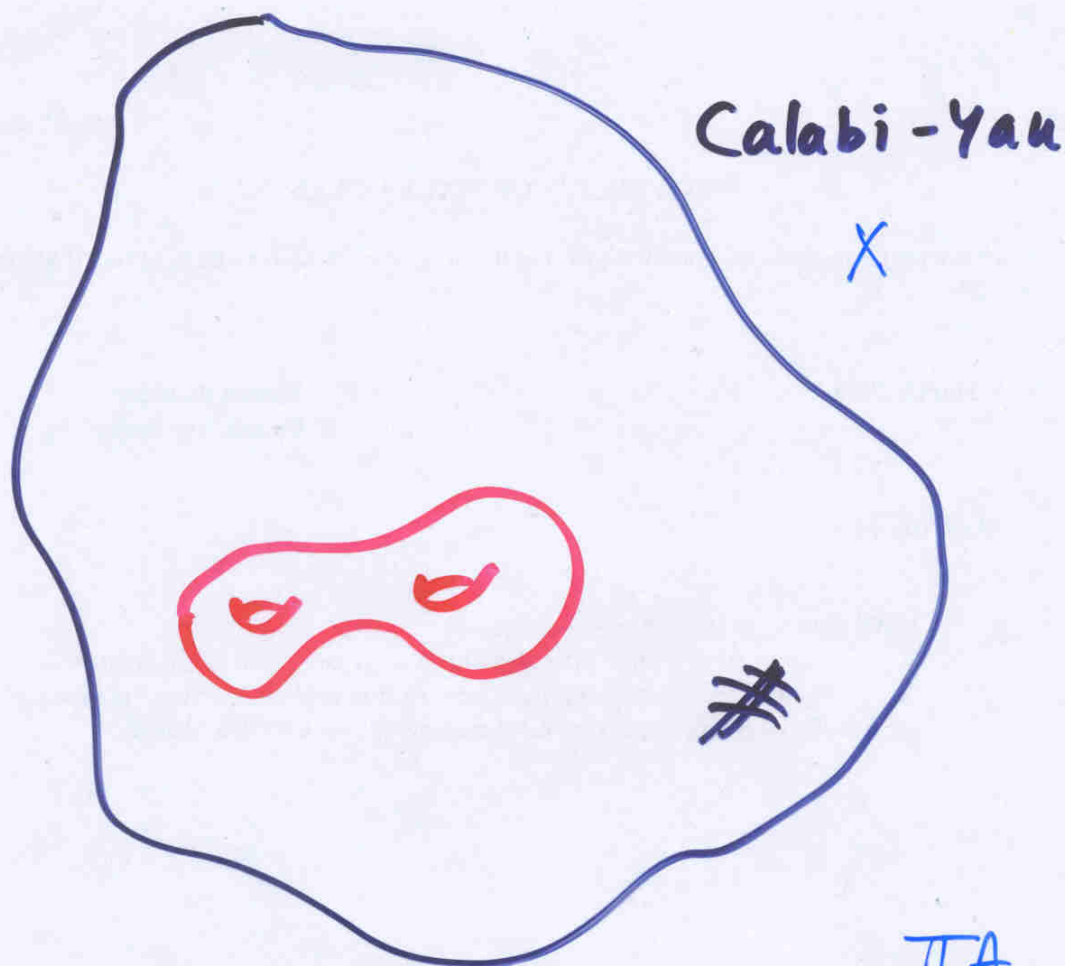
OF SUMMABLE STRING
PERTURBATION THEORY <sup>(beyond
2 dim)</sup>

PHYSICALLY INTERESTING

TYPE II SUPERSTRING

COMPACTIFICATION ON
CALABI-YAU SPACES

$\mathcal{N}=2$ FOUR DIMENSIONAL
SUSY



VECTOR MULTIPLETS - KÄHLER MODULI
 HYPER MULTIPLETS - COMPLEX MODULI

$$Z_g(a) = \sum_{\text{genus } g \text{ worldsheet } C \xrightarrow{f} X} \exp\left(-\int_C a\right)$$

GEOMETRY OF VECTOR MULTIPLETS -
 - WS INSTANTONS

$\mathcal{F}_0(a)$ - GENUS ZERO PART
 CALCULATES PREPOTENTIAL
 OF LOW-ENERGY EFFECTIVE
 THEORY

$\mathcal{F}_g(a)$ $g > 0$
 GRAVIPHOTON COUPLINGS

CAN WE CALCULATE
 ALL \mathcal{F}_g 's ?

WORLD SHEET POINT OF VIEW:

$$\mathcal{F}_g(a) = \sum_{\beta \in H_2(X, \mathbb{Z})} e^{-\int_{\beta} \omega_a} N_{\beta, g}$$

of holom maps:

$\Sigma_g \rightarrow X$
 $[\Sigma_g] = \beta$

TORIC CALABI-YAU MANIFOLDS

HAVE ENOUGH SYMMETRIES

CONSTRUCTION : linear σ -model
gauged susy

START WITH VECTOR SPACE

\mathbb{C}^{N+3} , coordinates

Z^A

$A = 1, \dots, N+3$

QUOTIENT BY THE TORUS

$$\mathbb{T}_{\mathbb{C}} = (\mathbb{C}^{\times})^N$$

$$Z^A \mapsto (\exp i Q_a^A \varphi^a) Z^A$$

$a = 1, \dots, N$

Q - $(N+3) \times N$
 \mathbb{Z} -valued
matrix

D-terms constraints :

$$\sum_A Q_a^A |Z^A|^2 = r_a = \text{const}$$

modulo $Z^A \mapsto \exp i(Q \cdot \varphi)^A Z^A$

$$\varphi^a \in \mathbb{R}/2\pi\mathbb{Z}$$

i.e. $T = U(1)^N$

THE RESULTING SPACE

$$X_3 = \mathbb{C}^{N+3} / (\mathbb{C}^*)^N$$

ADMITS

CALABI-YAU (RICCI FLAT KÄHLER)

METRIC iff

$$\sum_A Q_a^A = 0 \quad a=1, \dots, N$$

GEOMETRIC IMAGE OF TORIC MANIFOLD

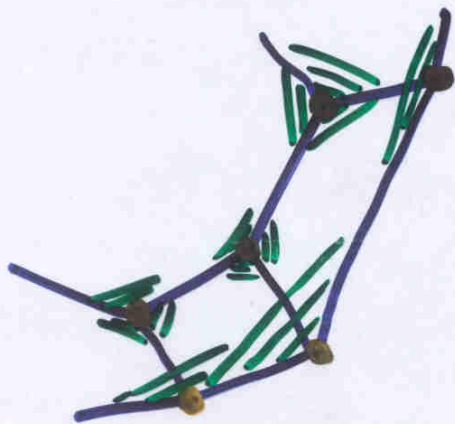
NEWTON POLYHEDRON

$$\Delta_X = X_3 / U(1)^3 \quad (\text{CONVEX})$$

$$\Delta_X \subset \mathbb{R}_{\geq 0}^{N+3}$$

"

$$\left\{ (P_1, \dots, P_{N+3}) \mid \sum_A Q_A^A P_A = r_a, \quad P_A \geq 0 \right\}$$



ORBITS OF T^3

OVER INTERIOR
POINTS OF $\Delta_X -$

$U(1)^3$

OVER FACES $U(1)^2$

EDGES $U(1)$

VERTICES \bullet

EXAMPLES :

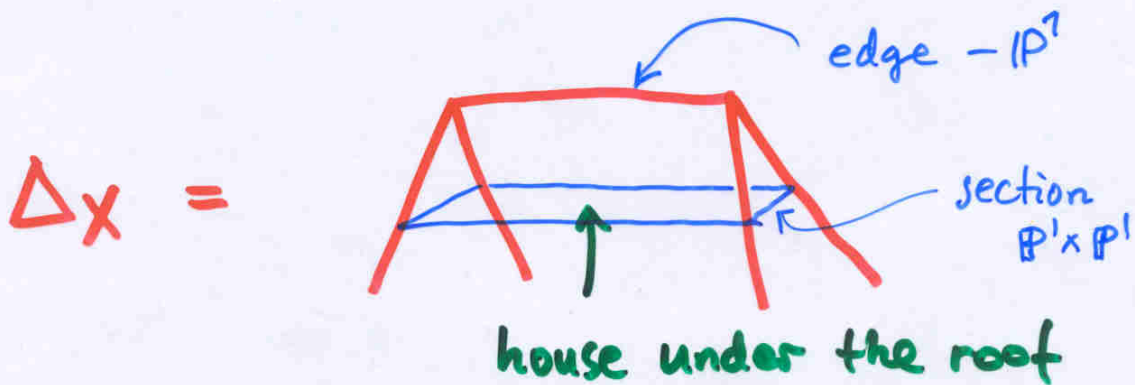
$$X_3 = \mathbb{C}^3$$

$$\Delta X =$$



$$\mathbb{R}_{30}^3$$

$$X_3 - \text{local } \mathbb{P}^1 \cong \begin{matrix} \mathcal{O}(-1) \oplus \mathcal{O}(-1) \\ \downarrow \\ \mathbb{P}^1 \end{matrix}$$



$$p_1 + p_2 - p_3 - p_4 = r, \quad p_A \geq 0$$

TOPOLOGICAL STRING OF TYPE A ON X_3

$$\begin{aligned}
 S = & \int p_{i\bar{z}} \partial_{\bar{z}} X^i + p_{\bar{i}\bar{z}} \partial_{\bar{z}} X^{\bar{i}} + \\
 & + \chi_{i\bar{z}} \nabla_{\bar{z}} \psi^i + \chi_{\bar{i}\bar{z}} \nabla_{\bar{z}} \psi^{\bar{i}} + \\
 & \alpha' (G^{i\bar{i}} p_{i\bar{z}} p_{\bar{i}\bar{z}} + R_{\mu\nu}{}^{i\bar{j}} \psi^\mu \psi^\nu \chi_{i\bar{z}} \chi_{\bar{j}\bar{z}}) \\
 & + \chi \bar{\chi} \partial \nu + \\
 & + \chi \bar{\chi} \Gamma_{\mu\bar{j}}^i G^{\bar{j}\bar{j}} V^\mu(\varepsilon) + \text{c.c.} \\
 & + \frac{1}{2\alpha'} (G_{\mu\nu} V^\mu(\varepsilon) V^\nu(\varepsilon) \\
 & + \psi \partial \psi \dots)
 \end{aligned}$$

$$\varepsilon \in \text{Lie } \mathbb{T}_\mathbb{C}^3$$

ω

$\bar{\varepsilon}$ - auxiliary (nothing depends on $\bar{\varepsilon}$, or α')

IN THE LINEAR SIGMA MODEL APPROACH

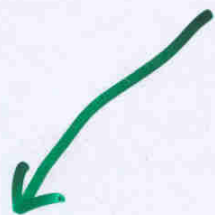
Σ - TWISTED MASSES

RESULTS :

$$X_3 = \mathbb{C}^3$$

$$Z(t, \varepsilon_1, \varepsilon_2, \varepsilon_3) = M(q) \frac{(\varepsilon_1 + \varepsilon_2)(\varepsilon_1 + \varepsilon_3)(\varepsilon_2 + \varepsilon_3)}{\varepsilon_1 \varepsilon_2 \varepsilon_3}$$

$$= \exp \sum_{g=0}^{\infty} t^{2g-2} F_g(\varepsilon)$$



$$\frac{1}{\varepsilon_1 \varepsilon_2 \varepsilon_3} \int_{\mathcal{M}_g} c_H(\varepsilon_1) c_H(\varepsilon_2) c_H(\varepsilon_3)$$

↑ Hodge bundle Chern polynomial
 $H^0(\Omega^1_{\Sigma})$

Here $M(q) = \prod_{n=1}^{\infty} \frac{1}{(1-q^n)^2}$

$$q = -e^{\frac{i\hbar}{t}}$$

$$M(q) \sim \exp\left(\frac{\zeta(3)}{t^2} + \sum_{g>1} \frac{|B_{2g}| |B_{2g-2}| t^{2g-2}}{2g(2g-2)(2g-2)!}\right)$$

asymptotic series
from top string

$M(q)$ - MacMahon function

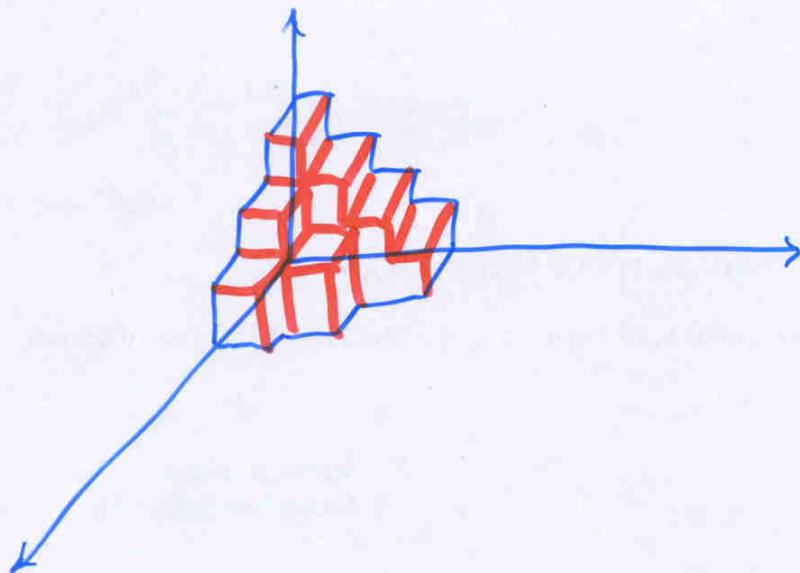
"
 \sum
 3d partitions
 R:

$$q^{|\pi|} \quad \text{also} =$$

$$= \exp \sum_{k=1}^{\infty} k q^k \mu(k)$$

↙ M. Green's function

" $\sum \frac{1}{d^k}$



Why 3d partitions?

$\pi \leftrightarrow \mathcal{J}_\pi$ - monomial ideals in $\mathbb{C}[z_1, z_2, z_3]$

$$\{ z_1^{i-1} z_2^{j-1} z_3^{k-1} \mid (i, j, k) \notin \pi \}$$

Why ideals? Why monomial? Why in \mathbb{C}^3 ?

see
below

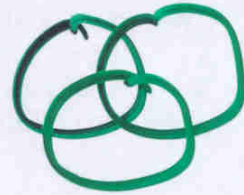
invariance
under
 \mathbb{T}^3

STRING ON
 \mathbb{C}^3

WE WANT ^{to} COUNT HOLOMORPHIC
CURVES IN X_3

CURVES ARE DESCRIBED BY
HOLOMORPHIC FUNCTIONS
(LOCALLY, POLYNOMIALS)

$$X_3 = \bigcup_{\alpha} U_{\alpha}$$



$\Sigma \subset X_3$ - holomorphic curve

$$\mathcal{I}_{\Sigma} \subset \mathcal{O}_{\Sigma} = \mathbb{C}[U_{\alpha}]$$



all polynomials in U_{α} , which vanish
on $\Sigma \cap U_{\alpha}$

COMPACT CURVES IN

\mathbb{C}^3 — POINT-LIKE

TOGETHER WITH INVARIANCE UNDER

TORS (GAUGING) THIS

ESTABLISHES 1-1 CORRESPONDENCE

WITH 3D PARTITIONS

WHY SHOULD WE COUNT

THEM SO SIMPLY

— ANOTHER STORY

(CALABI-YAU
CONDITION
IS IMPORTANT)

In some sense, toric X is built out of several \mathbb{C}^3 's
 one copy per fixed point of

$T^3 = U(1)^3$ action on $X \Leftrightarrow$
 vertex of Δ_X

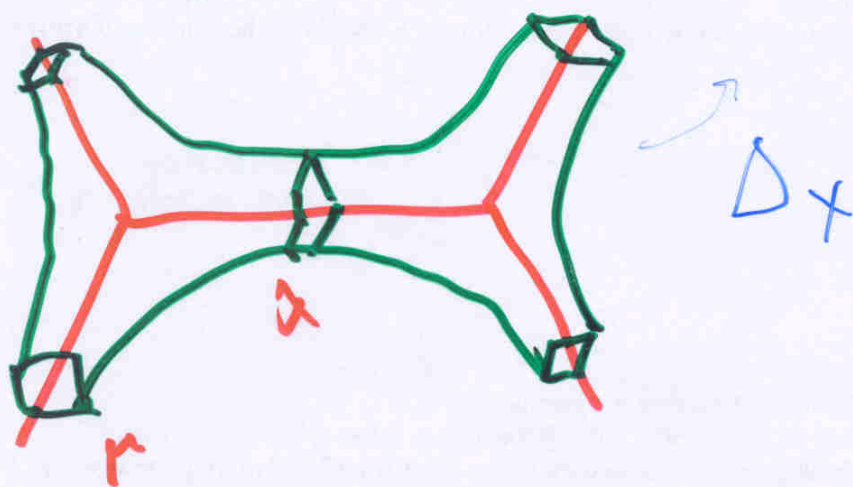
Tr_q
 Holomorphic functions χ

$$\text{Ch}_X = \sum_{\substack{\vec{n} \in \mathbb{Z}^{N+3} \\ \vec{Q} \cdot \vec{n} = 0}} q^{\vec{n}} = \leftarrow \text{twisted Witten index of SQM on } X$$

$$= \sum_{v \in \text{Vert } \Delta_X} \frac{1}{(1-q_{1v})(1-q_{2v})(1-q_{3v})}$$

$q_1, q_2, q_3|_v$ - local weights of T^3 action on $T_v X$

GW theory of toric X



also -
PANTS
DECOMPOSITION
FOR
WFT
(ask Robert
Dijkgraaf)

AT EACH VERTEX - \mathbb{C}^3 story
but, with non-trivial asymptotics
along the edges

2d partition

Kähler moduli

$$Z_X(\hbar) [t_e] = \sum_{\pi, \lambda_e} \pm q^{|\pi|} e^{-t_e |\lambda_e|}$$

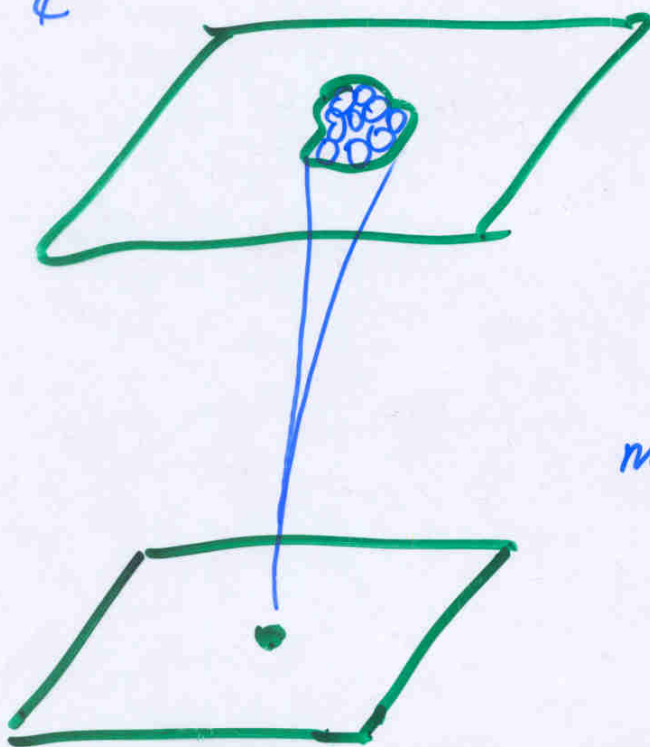
"CRYSTALS" (also: fluctuating geometry)

GO BACK TO \mathbb{C}^3 EXAMPLE

WE NOW GIVE ANOTHER
INTERPRETATION TO THE
CRYSTAL PICTURE

ALL
CONSIDER TORIC MANIFOLDS

X_3
 \downarrow
 \mathbb{C}^3 WHICH PROJECT TO



1-1 EXCEPT
AT 0

(BLOWUPS AT
 \nearrow
multiple 0)

NOW, PICK UP
THE PIECES

TOPOLOGICAL
STRING

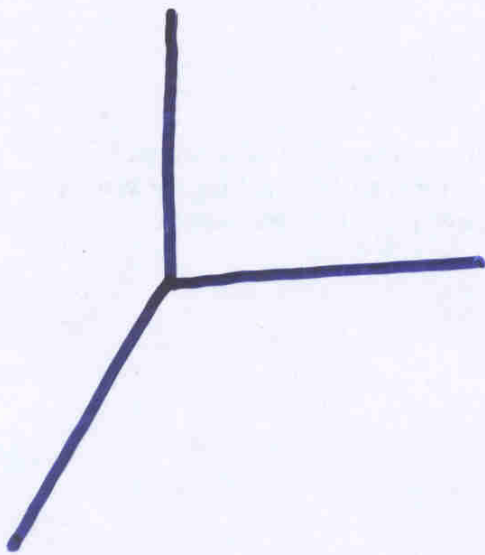
ON \mathbb{C}^3 (LARGE VOLUME LIMIT
OF ANY CALABI-YAU)

IDEALS OF POINTS
ON \mathbb{C}^3

TARGET
SPACE
THEORY

(TORIC) KÄHLER
GEOMETRIES OBTAINED
BY BLOWUPS
(FOAM)

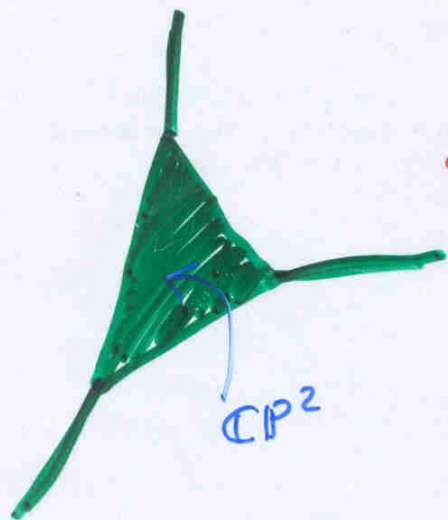
R^3



\mathbb{C}^3

Δ_x is obtained from

$\mathbb{R}_{>0}^3$ by cutting vertices

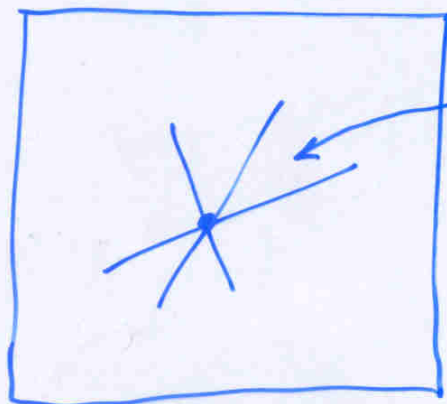


single blowup $\widetilde{\mathbb{C}^2}$

$$(|z^1|^2 + |z^2|^2 + |z^3|^2)^2 =$$

$$= r + |z^4|^2$$

$$r > 0$$

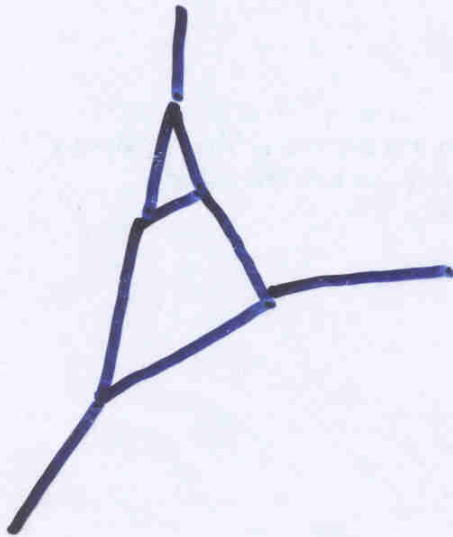


$\mathbb{C}P^2$ - space of lines

$$\widetilde{\mathbb{C}^2} = \{ (l, x) \mid x \in \mathbb{C}^3 \}$$

l - line, passing through 0 and x

DOUBLE BLOWUPS



CONTAINS NON-CONTR.
2-SPHERES

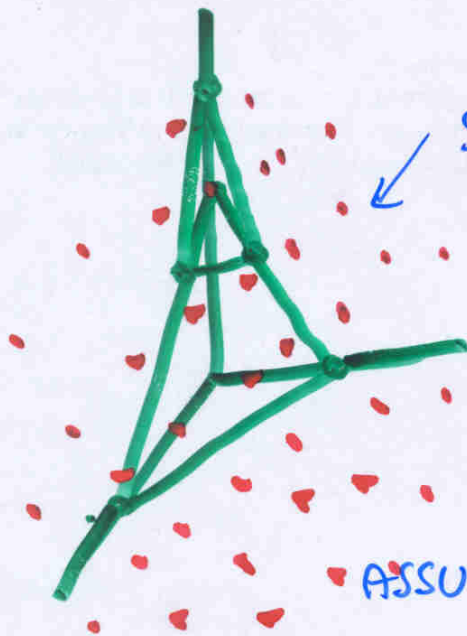
CONTINUE

FROM BLOWN UP GEOMETRY



3D PARTITIONS

GEOMETRIC QUANTIZATION



SOME TORIC SPACE

X_3 ,

IT HAS KÄHLER
FORM ω

ASSUME $[\omega] \in h H^2(X, \mathbb{Z})$

THEN QUANTUM STATES
FORM THE IDEAL



INTEGRAL
POINTS INSIDE

ΔX

THUS WE LEARN :

Z^A_X PARTITION FUNCTION OF TYPE A TOP. STRING
 $= \exp \sum_g \hbar^{2g-2} Z_g(a)$
IIA BPS

||

SUM OVER 3D CRYSTALS, EMBEDDED

INTO Δ_X

??

SUM OVER RANDOM

TORIC "GEOMETRIES"

??

LOCALIZED PATH INTEGRAL
IN SOME GRAVITY THEORY ON

X

another picture!
"Guinness"

MORE 3d PARTITIONS
THEN SMOOTH TORIC'S

IS THIS PICTURE REASONABLE?

WHAT SORT OF GRAVITY THEORY

DO WE EXPECT TO LIVE IN THE

TARGET SPACE OF TOPOLOGICAL

CLOSED STRING?

B MODEL IS SIMPLER

(CARES ABOUT COMPLEX
STRUCTURE ONLY)



Kodaira - Spencer F.T.

$$\int A' \frac{1}{\partial} \bar{\partial} A' + \int A' \wedge A' \wedge A' \Omega^2$$

A MODEL IS TRICKIER (KÄHLER STRUCTURES)

$$\int \text{KIN. TERM} \quad + \quad \int k \wedge k \wedge k + \dots$$

$k \frac{1}{d^c} d k$



IN OUR COUNTING WE SEE

$$k = k_0 + \hbar F \quad \leftarrow \begin{array}{l} \text{quantized} \\ \in H^2(X, \mathbb{Z}) \end{array}$$

↑
ORIGINAL
KÄHLER FORM
on X_3

↑
LOOKS LIKE
DISCRETIZATION
(cf. 2d MATRIX
MODELS)

EACH CRYSTAL

(mathematically they correspond to the so-called torsion free sheaves of rank 1)

is characterised by :

$$\mathcal{O}_X \sim F \wedge F$$

$$\mathcal{O}_X \sim F \wedge F \wedge F$$

CRYSTAL FORMULA

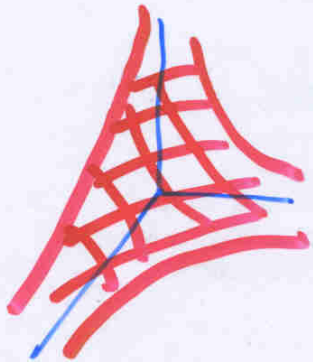
$$\sum_{\text{crystal}} \exp \left(\int_X k_0 \wedge \mathcal{O}_X + t_1 \int_X \mathcal{O}_X \right)$$

\downarrow
 $\int k^3$

REMARKS

1. $\hbar \rightarrow 0$ CLASSICAL GEOMETRY

(NOT THE ORIGINAL ONE)



LIMIT SHAPE -
- MIRROR OF X_3

2. IN PARTICULAR, FOR

$$X_3 = \widetilde{\mathbb{C}^2/\Gamma}$$

\downarrow
 P_1

IN 4d $N=2$ SYM

WITH GAUGE GROUP
A, D, E

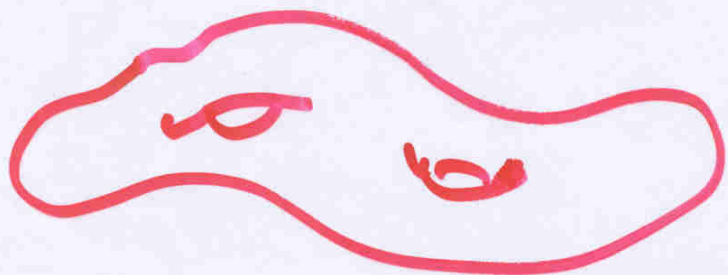
SW
CURVES

CRYSTAL
SUM

||
4d INST
SUM

3. S-duality

VERY DIFFERENT PACKAGING
OF INSTANTON + PERT.
SUM



holomorphic
curve
in X_3

A model
perturbation
theory

β model-
like
D-instanton
sum

$$\hbar = g_A^s = \frac{1}{g_B^s}$$

$$\sum_{g,m} e^{-m \text{Area}(C)} N_{g,m} \hbar^{2g-2} = \log \left(\sum_J e^{-\hbar k^3} N_J e^{-\int C h_2 J \wedge k_g} \right)$$

5. Duality continues to hold in more complicated backgrounds (non-CY)



SAY, $U(N+1) \times U(N)$ gauge theory $N=2$

AT SPECIAL POINT ON THE COULOMB BRANCH

(RELATED TO B model

on $P^1 \times C^2$, S-dual to A model + top-gravity sector turned on)

$$Tr_1 \phi^k - Tr_2 \phi^k \leftrightarrow \sigma_{k-1}(\omega)$$

CONCLUSIONS

- WE SEE GLIMPSES OF FLUCTUATING SPACE-TIME GEOMETRY (FOAM) IN THE FULL ALL-GENUS TOPOLOGICAL STRING PARTITION FUNCTION
 - SO FAR ON TORIC MANIFOLDS (BUT THE CONJECTURE MAKES SENSE ON ANY CALABI-YAU)
- S-duality
- Kodaira-Spencer + self-dual tensor ?
- M-TOPOLOGICAL THEORY ?
 G_2 ? E_{10} ? $\mathbb{C}P^{3/4}$?