

# Electric-magnetic duality and all that

W. Nahm, DIAS

diofest

- 
- a simple duality
- two papers and a talk
- $M + \bar{\Phi}$
- $E$ -diagrams

$S$ -matrix  $\longleftrightarrow$  CFT

Consider an integrable QFT<sub>2</sub> on a circle of circumference  $L$ .

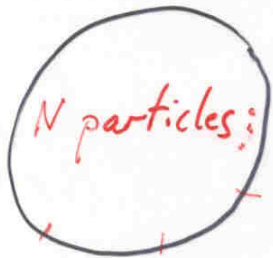
S-matrix  $\xrightarrow{L} 0$  CFT

Bethe ansatz for the wavefunction of two particles of rapidities  $\theta_1, \theta_2$ .



$$\psi(x_1, x_2) = \begin{cases} e^{ik_1 x_1 + ik_2 x_2} & \text{for } x_1 \ll x_2 \\ e^{i\delta(\theta_1 - \theta_2)} e^{ik_1 x_1 + ik_2 x_2} & \text{for } x_1 \gg x_2 \end{cases}$$

S-matrix for  $\theta_1 > \theta_2$



$$\sum_{j \neq i} \delta(\theta_i - \theta_j) + k_i L = 2\pi n_i \quad i = 1, \dots, N$$

For  $L \rightarrow 0$ , left- and right-movers decouple.

Energy of right movers  $\approx$  momentum of right movers.

Momenta are quantized and easy to calculate.  
 $\Rightarrow$  CFT energies are easy to calculate ?!

Two particles:

$$\delta(\theta_1 - \theta_2) + k_1 L = 2\pi u_1$$

$$\delta(\theta_2 - \theta_1) + k_2 L = 2\pi u_2$$

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$$\delta(\theta_1 - \theta_2) + \delta(\theta_2 - \theta_1) = -2\pi A \quad A \in \mathbb{Z}$$

$N$  particles:

$$\sum_{\substack{i,j=1 \\ i \neq j}}^N \delta(\theta_i - \theta_j) + L \sum_i k_i = 2\pi \sum_i u_i$$

$$L \sum k_i = 2\pi \sum u_i + 2\pi \binom{N}{2} A$$

assume Bose statistics,  $n_i \geq 0$

$$\text{tr} \rho q^{L_0} = \sum_{N=0}^{\infty} \frac{q^{\binom{N}{2} A + 6N + 4}}{(q)_N} \Rightarrow \text{central charge} \dots$$

Application for non-unitary models:

Let  $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-$ ,  $\mathcal{H}_{\pm}$  graded by  $L_0$ ,  
metric positive definite on  $\mathcal{H}_+$ , negative on  $\mathcal{H}_-$ .

For single particle states in  $\mathcal{H}_-$  (as in the  
(2,5) minimal model)

$$\text{tr}_{\mathcal{H}_+} q^{L_0} - \text{tr}_{\mathcal{H}_-} q^{L_0} = \sum_{N=0}^{\infty} (-1)^N \frac{q^{\binom{N}{2} A + 6N + 4}}{(q)_N}$$

(2,5) model:  $A=2$

# Gauge theory and magnetic charge

P. Goddard, J. Nuyts, D.I. Olive, Dec. 1976

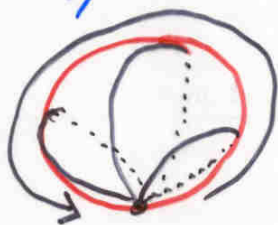
gauge field, Higgs field  $\varphi$   
 gauge group  $G$  broken to  $H \subset G$

$\varphi(S^2_\infty)$  in single orbit of  $G$   
 $G/H$

$$\pi_2(G) \rightarrow \pi_2(G/H) \rightarrow \pi_1(H) \rightarrow \pi_1(G)$$

$\overset{H}{\circlearrowleft}$

$\pi_2(G/H) \cong \pi_1(H)G$   
 Higgs topology      gauge field topology



Replace  $H$  by maximal torus  $T = \mathbb{R}^r/\Lambda$

$\pi_1(T) = \Lambda$  monopoles (modulo Weyl group)

Representations of  $H \Rightarrow$  repr. of  $T$

$\Lambda^\vee$  electric charges (modulo Weyl group)

Conjecture: There is a magnetic gauge group  $H^\vee$  with weight lattice  $\Lambda$ .

Kernel of  $\pi_1(T) \rightarrow \pi_1(H)$ : unstable monopoles.

Preference for representations for which the weights form a single orbit of the Weyl group (minimal representations).

Words in quotation marks:

"Weyl group", "Weyl reflection"  
"section" (ref. Wu, Yang)  
"maximal torus"  
"centre"  
"universal covering group"  
"weight", "root"

Not in quotation marks:

Cartan metric, Killing form

Speculations:

• True symmetry group is  $H \otimes H^\vee$   
• analogy to Thirring model ~ Sine-Gordon theory  
topological charge, current  $\longleftrightarrow$  Noether charge, current

• monopoles ~ quarks (confined)

$H = SU(3)$  colour

$H^\vee = SU(3)/Z_3$  electroweak  
(citation H. Fritsch + P. Minkowski)

# Magnetic monopoles as gauge particles

C. Montonen, D. Olive, Oct. 1977

- $gg \in 2\pi\hbar\mathbb{Z}$  (Dirac)

$$(g, g) \rightarrow (g, -g)$$

- BPS mass formula  $M \sim \sqrt{g^2 + g^2}$

- forces in BPS limit (Montonen)

- moduli space:

charge lattices /  $SO(2)$

$$[\text{now: } SL(2, \mathbb{Z}) \backslash SL(2, \mathbb{R}) / SO(2)]$$

Speculation: Spin =  $1\hbar$ .

Higgs: Goldstone boson of  $SO(2) \subset SU(2, \mathbb{R})$ .

E. Witten and D. Olive (1978):

BPS mass formula as consequence of supersymmetry.

Mass  $\sim$  central charge

# Magnetic monopoles and electromagnetic duality conjectures.

D. Olive, talk at Monopole Meeting, Trieste Dec. 1981  
(no citations)

• are there local monopole fields?  
(cf.  $\phi^4$  solitons versus  $\cos\phi$  solitons)

• susy  $N=2$  or  $N=4 \Rightarrow \text{spin } 1/2$

"The [fermionic] zero modes .. indicate what sort of multiplet the monopoles lie in. The fermionic zero modes are represented by Dirac gamma matrices indicative of the varying spin content."

• renormalization of  $q, g$

$$q(\Lambda) \cdot g(\Lambda) \in 2\pi\hbar\mathbb{Z}$$

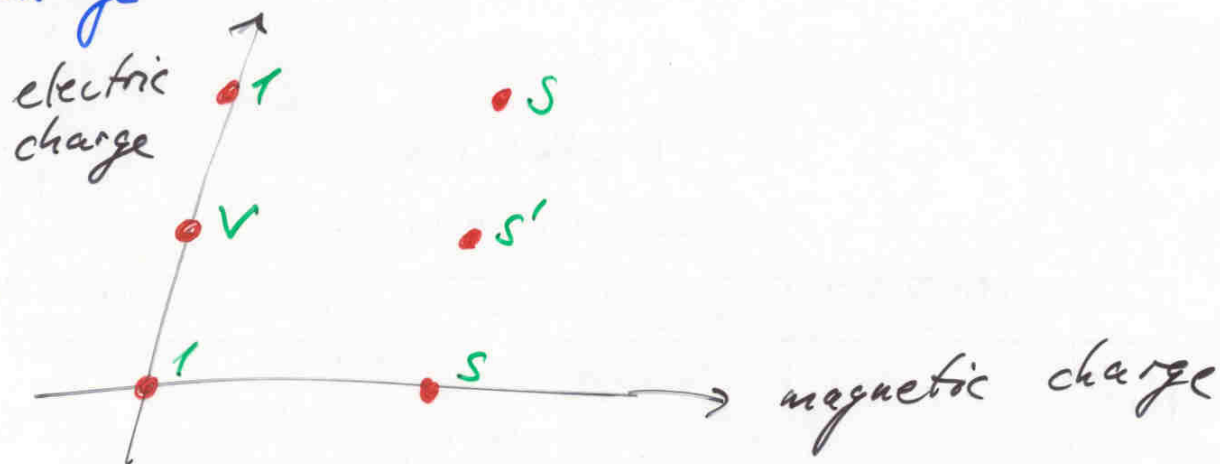
$$q(\Lambda) \sim g(\Lambda)$$

$\Rightarrow$  no renormalization,  
 $\beta = 0$

"... the  $N=4$  supersymmetric extension .. seems to be the only theory for which the Montonen Olive conjecture could be valid and .. its validity is reasonably plausible at least for  $G = SU(2)$ ."

Why did we overlook  $N_f = 4$ ?

charge lattice for  $N_f = 4$  (massless)



$SO(2N_f)$  representations:

$v$  vector

$s, s'$  half-spinors

Full  $SL(2, \mathbb{Z})$  electric-magnetic duality by using  $SO(8)$  triality.

For  $N_f = 0, 2, 3$ :

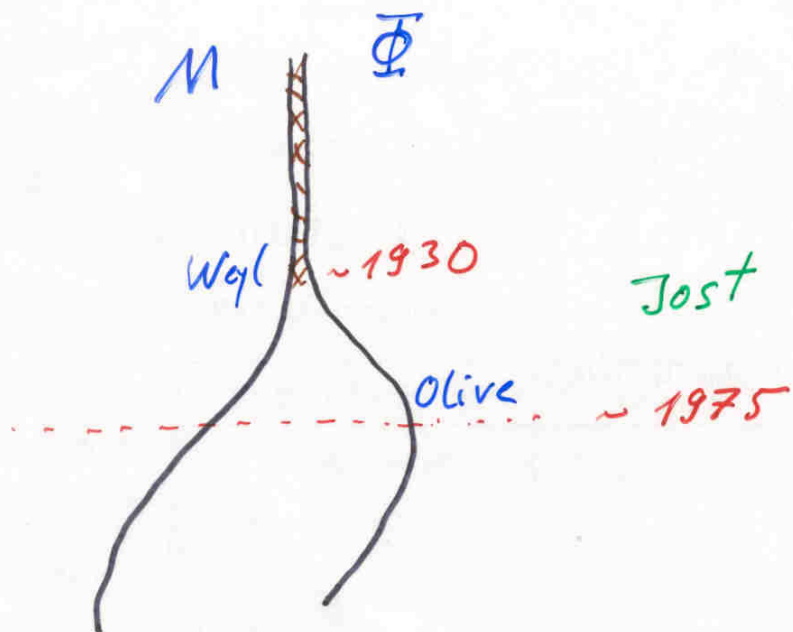
finite index subgroup of  $SL(2, \mathbb{Z})$

$N_f = 1$ : Argyres-Douglas singularity

in moduli space  $\Rightarrow$  double cover.

"in the  $N=2$  theory the monopoles lie in a multiplet not containing spin  $t$ ."





QFT + geometry

solitons  
19th century

Skyrme

Coleman

+ Hooft, Polyakov

gauge theory

Weyl

Dirac

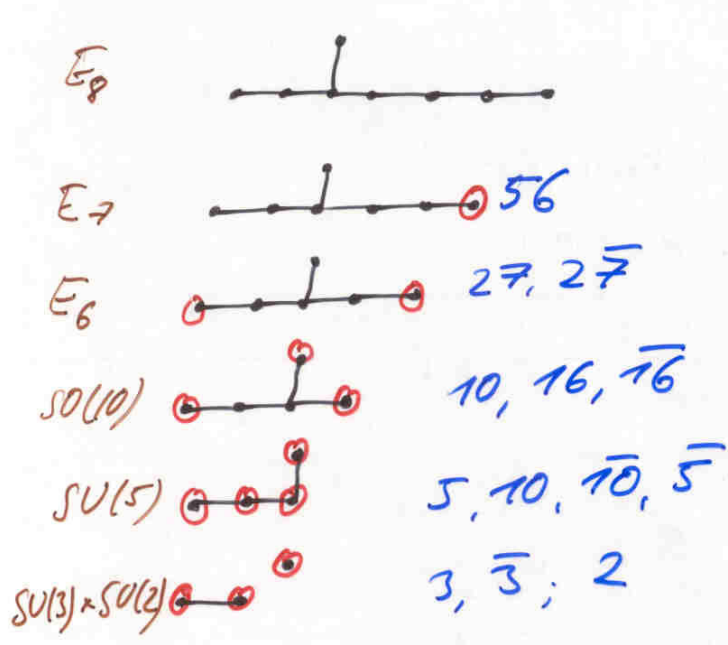
Yang

susy

GSO  
Witten, Olive

'76 Goddard, Nuyts, Olive  
'77 Montonen, " duality

'94 Sen  
Seiberg, Witten

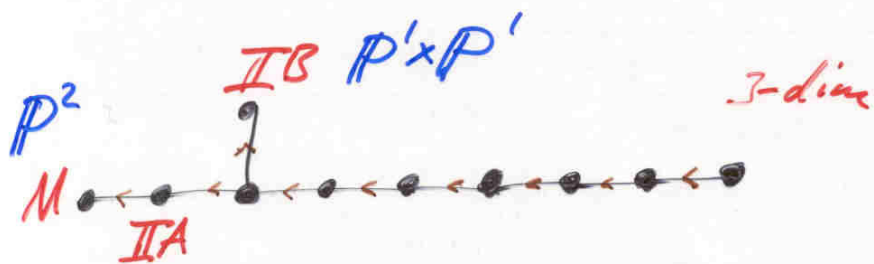
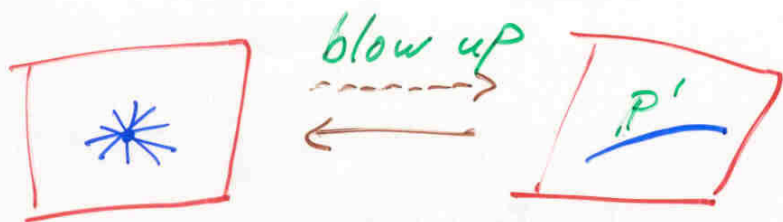


"the groups  $E_8$  to  $E_3$  (with possible  $U(1)$  factors) are candidates for grand unified groups of the fundamental interactions with the fundamental fermions assigned to the minimal representation."

"Mathematically there is nothing very special about the groups  $SU(3) \times SU(2)$ ,  $SU(5)$  or  $SO(10)$  per se, but they do have a special significance as the lowest members of the sequence of  $E_n$  groups."

An E-series also appears in the torus compactification of 11-dim supergravity (E. Cremmer, B. Julia). The  $SL(2, \mathbb{R})$  electric-magnetic duality group of the sourceless Maxwell equations can arise in this way.

$E_{n-d}$  for compactification to  $d$  dimensions. Relation to Del Pezzo surfaces (Lepol, Neitzke, Vafa).



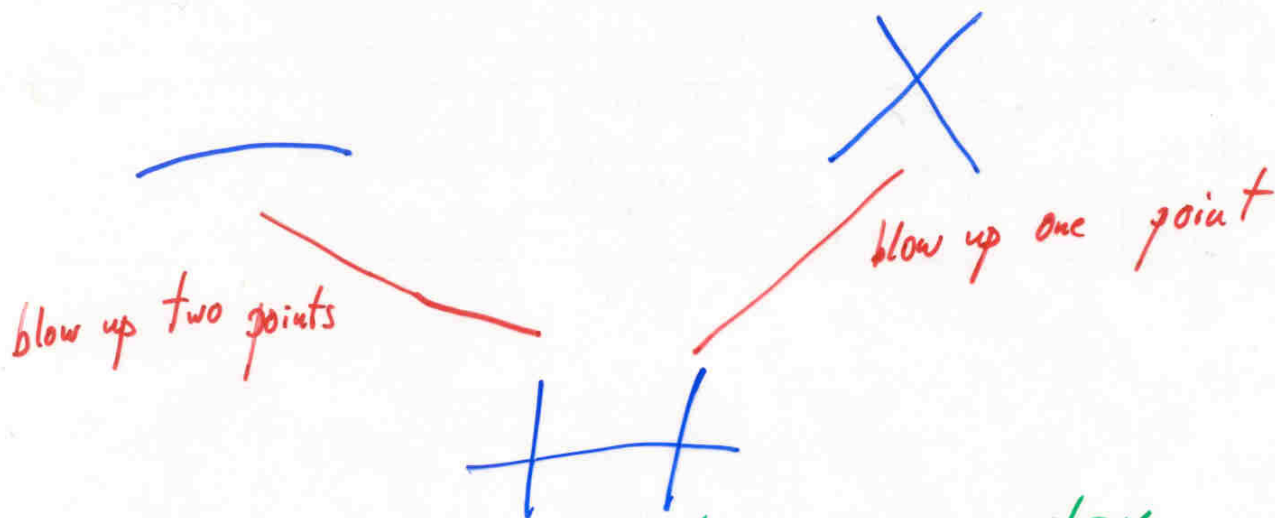
10 Del Pezzo surfaces

$\mathbb{P}^2$ : Affine  $\mathbb{C}^2 + \mathbb{P}^1$  at  $\infty$ .

$\mathbb{P}^1 \times \mathbb{P}^1$ :  $xy = uv \quad (x, y, u, v) \in \mathbb{P}^3$

For  $x=1$ : affine  $\mathbb{C}^2 \ni \{u, v\}$

$x=0$ : two lines ( $u=0$  or  $v=0$ ).  $\times$



Most elementary algebraic geometry yields E-diagrams.