

LATTICE QCD AT FINITE TEMPERATURE: SOME ASPECTS RELATED TO CHIRAL SYMMETRY*

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*Received 8 September 2022, accepted 10 October 2022,
published online 14 December 2022*

Out of the many exciting results obtained with the lattice approach to QCD under extreme conditions, I discuss a few selected items related to chiral symmetry: the chiral condensate as an approximate order parameter, meson screening masses, and masses of baryons and mesons, including $D_{(s)}$ mesons, when approaching the crossover from the hadronic side.

DOI:10.5506/APhysPolBSupp.16.1-A12

1. Introduction

Lattice QCD is *the* first-principle method to study QCD under extreme conditions, at least along the temperature axis in the QCD phase diagram and while considering equilibrium conditions. Indeed, studies of QCD thermodynamics, including *e.g.* the equation of state and fluctuations of conserved charges, are nowadays routinely done at the physical point while taking the continuum limit [1–4]. Extensions to nonzero baryon density have to deal with the sign problem, which can be confidently handled either at real but small quark chemical potential, or at imaginary quark chemical potential, after analytical continuation.

A clear next step, going beyond thermodynamics and the phase structure, is to consider the behaviour of hadrons. Heavy quarks (charm and bottom) and quarkonium states (charmonium and bottomonium) have been in the spotlight for a long time, due to their relevance for heavy-ion phenomenology. In this presentation, I will not be able to do justice to these developments. Instead, I will focus on the light-quark sector, emphasizing

* Presented at the 29th International Conference on Ultrarelativistic Nucleus–Nucleus Collisions: Quark Matter 2022, Kraków, Poland, 4–10 April, 2022.

chiral symmetry and its impact beyond the identification of the transition. In particular, I will discuss three aspects: the chiral condensate as an approximate order parameter; mesonic screening masses; baryon and meson masses, especially when approaching the crossover from the hadronic side.

2. Chiral symmetry and the thermal crossover

Chiral symmetry is an approximate symmetry, spontaneously broken in the hadronic phase. Observables linked to chiral symmetry are used to identify the thermal crossover, notably the chiral condensate and its fluctuations, the chiral susceptibility χ (see Fig. 1). In fact, the most precise determination of the pseudocritical temperature has been obtained from the latter, with the recent values

$$\begin{aligned} T_{\text{pc}}^{\chi} &= 158.0(6) \text{ MeV} && \text{Budapest–Wuppertal–Houston [5],} \\ T_{\text{pc}}^{\chi} &= 156.5(1.5) \text{ MeV} && \text{HotQCD [6].} \end{aligned}$$

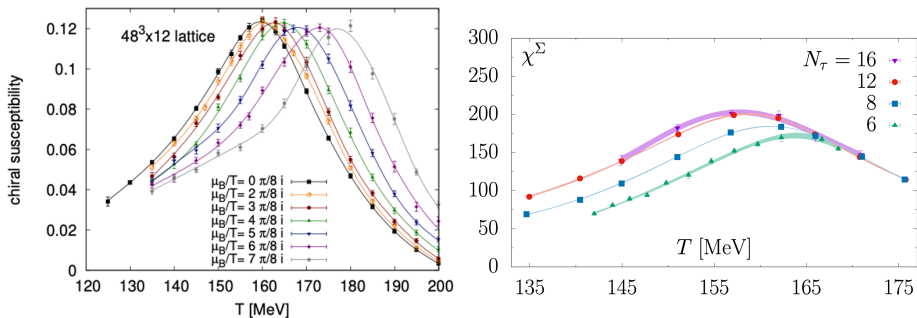


Fig. 1. Chiral condensate as a function of temperature, for several values of an (imaginary) baryon chemical potential (left, Budapest–Wuppertal–Houston [5]), and several values of the lattice spacing $a \sim 1/N_\tau$ (right, HotQCD [6]).

From a lattice perspective, the results quoted here have been obtained using staggered quarks, for which simulations are carried out at the physical point and a continuum extrapolation is feasible. The Wilson-type quarks are more expensive to simulate and hence results are not available at both the physical point and in the continuum limit. Yet there is steady progress here as well, as demonstrated in Fig. 2, which includes results from the twisted-mass [7] and Wilson-clover [8] fermions. These are not both at the physical point (*i.e.* $m_\pi > m_\pi^{\text{nature}}$) and in the continuum limit; nevertheless, an extrapolation in m_π using the Wilson-type quarks only yields $T_{\text{pc}}^{\psi\psi} = 158(3) \text{ MeV}$ at the physical point [8], in agreement with the numbers quoted above.

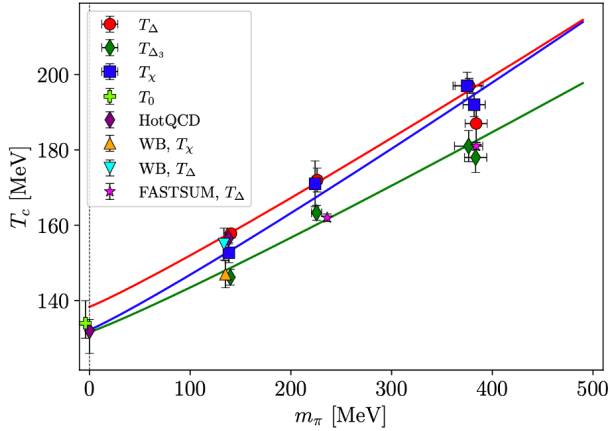
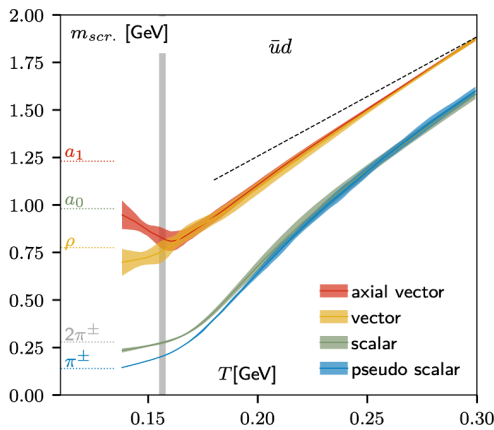


Fig. 2. Pion mass dependence of the pseudocritical temperature, obtained using the twisted-mass (first three symbols), staggered (HotQCD, WB), and Wilson-clover (FASTSUM) fermions, for several quantities related to the chiral susceptibility [7].

3. Mesonic screening masses

Going beyond thermodynamics, one may ask how chiral symmetry restoration affects hadrons and quarks. As a first step, we consider meson screening masses (or spatial screening lengths). These are computable on the same lattices as used for thermodynamics, since the computation benefits from large spatial volumes, while the Euclidean time is integrated over. Screening masses interpolate between the pole mass at $T = 0$ and



$T(\text{GeV})$	$\frac{m_P}{2\pi T}$	$\frac{m_V}{2\pi T}$
164.6(5.6)	1.0194(25)	1.0261(23)
82.3(2.8)	1.0219(15)	1.0291(18)
51.4(1.7)	1.0216(16)	1.0312(18)
32.8(1.0)	1.0217(15)	1.0302(19)
20.63(63)	1.0220(15)	1.0343(17)
12.77(37)	1.0185(18)	1.0306(24)
8.03(22)	1.0200(18)	1.0341(28)
4.91(13)	1.0192(18)	1.037(3)
3.040(78)	1.0124(18)	1.0380(25)
2.833(68)	1.0147(24)	1.038(3)
1.821(39)	1.0122(18)	1.044(4)
1.167(23)	1.0039(20)	1.045(6)

Fig. 3. Left: screening masses as a function of temperature, for $\bar{u}d$ mesons [9]. Right: screening masses normalised with $2\pi T$, up to *very* high temperature [10].

$\sqrt{(2\pi T)^2 + m_q^2} \approx 2\pi T$ at high T . They can, therefore, be used to probe both chiral symmetry restoration as well as the perturbative high-temperature limit.

This is demonstrated in Fig. 3: screening masses in the light-quark sector are shown on the left [9], with, in particular, the degeneracy in the vector (ρ) and axial-vector (a_1) channels emerging at or close to the transition. Normalised screening masses up to *very* high temperatures (note that units are GeV) are given on the right, including small but numerically significant deviations from $2\pi T$, which can be analysed in the perturbation theory [10].

4. Baryon and meson masses

Phenomenologically, the interest lies not so much in screening masses but rather in hadron masses. Ideally, information on the spectrum and in-medium modification is extracted from spectral functions $\rho(\omega, \mathbf{p}; T)$. Unfortunately, the inversion of the integral relating the Euclidean correlator and the spectral function is a classic ill-posed problem. The next best thing is to start with the assumption that hadrons exist at low temperature, with a computable mass for the ground-states, and subsequently, raise the temperature of the hadron gas to follow the response. Ground-state masses are extracted from temporal lattice correlators, which benefits from having access to many Euclidean time points. This motivates the use of anisotropic lattices, with $a_\tau \ll a_s$, leading to quite different lattice geometries compared to the thermodynamic studies.

In the following I will report on work by the FASTSUM Collaboration [8]. Briefly, FASTSUM uses $N_f = 2 + 1$ flavours of the Wilson-clover quarks, with an anisotropy of $a_s/a_\tau \sim 3.45$ and a lattice cutoff of $a_\tau^{-1} \sim 6$ GeV. Light quarks are heavier than in nature; results are shown for two sets of ensembles: Generation 2 with $m_\pi = 384(4)$ MeV and Generation 2L with $m_\pi = 236(2)$ MeV. In Gen2 (2L), we have 4(5) ensembles in the hadronic phase and 5(6) in the quark-gluon plasma. The transition temperature as determined by the chiral susceptibility is 170(4) and 165(3) MeV, respectively.

Let us start with baryons [11–13]. Chiral symmetry breaking is manifested here by the absence of parity doubling, *i.e.* in vacuum positive- and negative-parity states are nondegenerate¹. One can show that if chiral symmetry is unbroken, parity doubling is present, not only for states, but already at the level of Euclidean correlators [12]. Hence this can be investigated without any further assumptions on the spectral content. Denoting the positive/negative-parity correlator with $G_{+/-}$, we may construct a quasi-order parameter

¹ Consider *e.g.* the nucleons $m_N = m_+ = 939$ MeV and $m_{N^*} = m_- = 1535$ MeV.

$$R = \sum_{\tau} \frac{G_+(\tau) - G_-(\tau)}{G_+(\tau) + G_-(\tau)},$$

which interpolates between 1 at low temperature — assuming that the groundstates m_{\pm} dominate and that $m_- \gg m_+$ — and 0 when chiral symmetry is restored. This R parameter is shown in Fig. 4 for octet and decuplet baryons, for both sets of ensembles [8]. We observe that indeed R changes from about one to close to zero, with the inflection points shifting to lower temperature as the light quarks get lighter. The inflection-point temperatures for the ensembles with the lighter pion lie between 157 and 160 MeV, close to the thermal crossover.

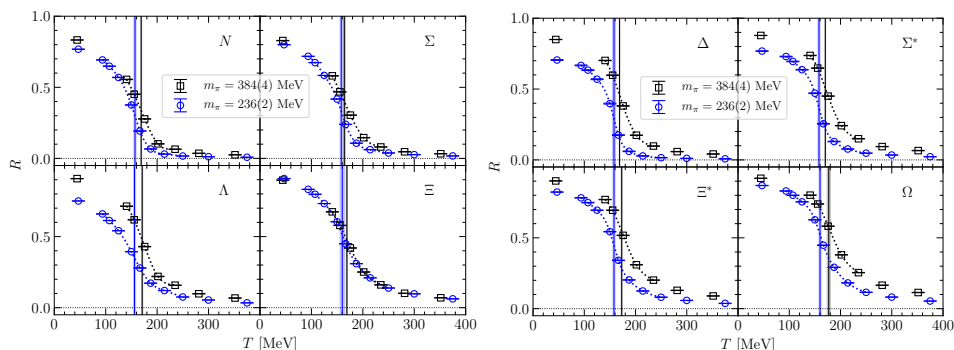


Fig. 4. Parity-doubling parameter R as a function of temperature for octet (left) and decuplet (right) baryons, for two sets of ensembles labelled by the pion mass. The vertical lines indicate the inflection points, which shift to lower temperature as the light quarks get lighter [8].

The analysis can be taken a step further, by extracting the ground-state masses from the lattice correlators, for both parities. The results for the ensembles with the heavier pion are shown in Fig. 5, in the hadronic phase. Note that they are normalised with m_+ at the lowest temperature. We observe that the positive-parity masses are mostly independent of temperature, while the negative-parity masses are reduced, such that an approximate degeneracy emerges at the thermal crossover (recall that $T_{\text{pc}} = 170$ MeV for these data). While the degeneracy is expected, the manner in which this is realised can only be found by an actual computation. These results can be used as a benchmark for effective parity-doublet models used at finite density, see *e.g.* Ref. [14].

Moving now to the light-meson sector, I present here some preliminary FASTSUM results obtained by assuming that spectral broadening can be ignored at low temperature. This assumption can be investigated by a detailed comparison of correlators at different temperatures [15]. Results for

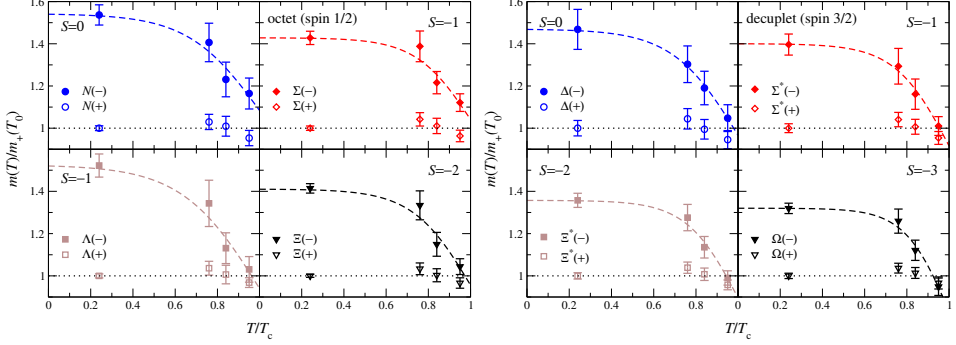


Fig. 5. Temperature dependence of the ground-state masses, normalised with m_+ at the lowest temperature, in the hadronic phase, for octet (left) and decuplet (right) baryons, for the Gen2 ensembles. Positive- (negative-) parity masses are indicated with open (closed) symbols [13].

the temperature dependence of the pion and kaon ground-states are shown in Fig. 6 (left). The light quarks are heavier than in nature, which mostly affects the pion. Note that the masses of both states appear to increase slightly as the thermal crossover is approached. On the right, the temperature dependence of the ρ (vector) and a_1 (axial-vector) ground-states is shown. The ρ state is approximately independent of temperature, while the a_1 state shows a rather strong temperature dependence. A degeneracy is a signal of $SU(2)_L \times SU(2)_R$ chiral symmetry restoration, which is seen to occur at or very close to the thermal transition, not dissimilar as in the case

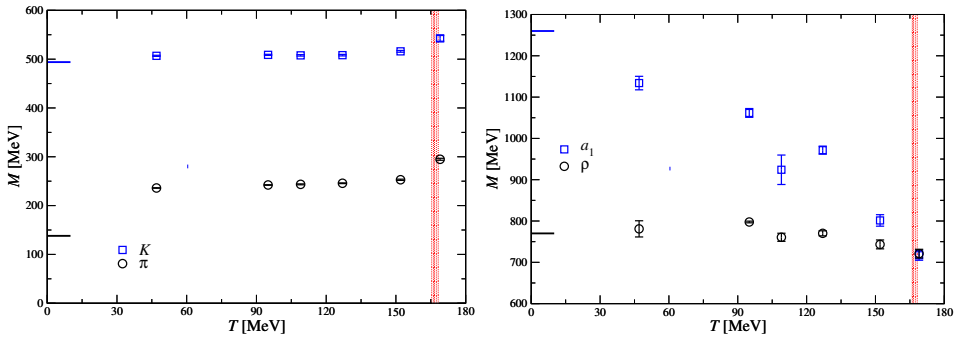


Fig. 6. Temperature dependence of the pion and kaon (left) and the ρ and a_1 (right) ground-state in the hadronic phase, for the Gen2L ensembles. The vertical band indicates the thermal transition, while the horizontal stubs at $T = 0$ represent the PDG value; note that the light quarks are heavier than in nature, which mostly affects the pion [15].

of screening masses, see Fig. 3 (left). We emphasise that these results are obtained under the assumption that narrow spectral functions can be used to describe the data, which might not be justified, especially in the axial-vector channel. It is noted that there are many model predictions for the thermal behaviour in these channels, see *e.g.* Ref. [16] and references therein.

Finally, we present some FASTSUM results for D and D_s mesons, in the pseudoscalar and vector channels, see Fig. 7. In this case, the masses are decreasing as the thermal transition is approached, quite similar, in fact, to the behaviour seen in effective models [17].

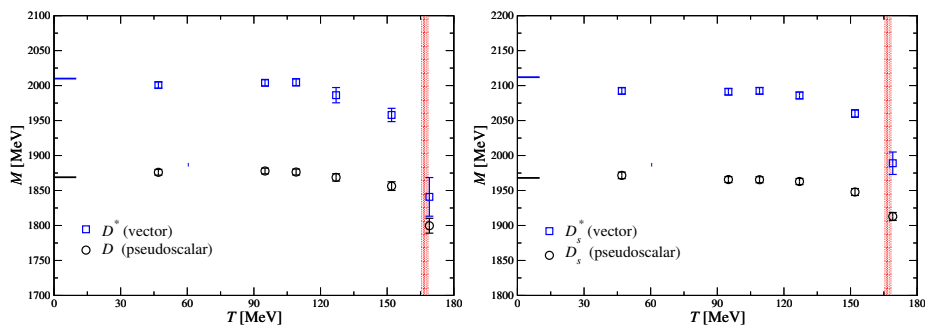


Fig. 7. As in Fig. 6 for the D and D^* (left) and the D_s and D_s^* (right) mesons [15].

5. Summary

As it is well known, chiral symmetry is of utmost importance for thermal QCD. In lattice studies, it is used to study properties of thermal crossover, providing the most precise estimates of the pseudocritical temperature. Going beyond thermodynamics, the expected restoration of chiral symmetry leads to emerging degeneracies in the spectrum, in the case of meson screening masses, for baryons and parity doubling, and for mesons and chiral partners. I have presented (new) lattice data in the hadronic phase, indicating a precursor to chiral symmetry restoration already in the hadronic phase. These results may be compared to effective model descriptions. This is especially relevant when the latter are extended to regions of the QCD phase diagram where lattice QCD is not directly applicable

The work presented in Sec. 4 was carried out by FASTSUM and I thank my colleagues for the fruitful collaboration. FASTSUM acknowledges DiRAC, PRACE, and Supercomputing Wales for the use of computing resources. I am supported by the UKRI Science and Technology Facilities Council (STFC) Consolidated Grant No. ST/T000813/1.

REFERENCES

- [1] S. Borsányi *et al.*, *J. High Energy Phys.* **2010**, 077 (2010), [arXiv:1007.2580 \[hep-lat\]](#).
- [2] HotQCD Collaboration (A. Bazavov *et al.*), *Phys. Rev. D* **90**, 094503 (2014), [arXiv:1407.6387 \[hep-lat\]](#).
- [3] S. Borsányi *et al.*, *J. High Energy Phys.* **2012**, 138 (2012), [arXiv:1112.4416 \[hep-lat\]](#).
- [4] HotQCD Collaboration (A. Bazavov *et al.*), *Phys. Rev. D* **86**, 034509 (2012), [arXiv:1203.0784 \[hep-lat\]](#).
- [5] S. Borsányi *et al.*, *Phys. Rev. Lett.* **125**, 052001 (2020), [arXiv:2002.02821 \[hep-lat\]](#).
- [6] HotQCD Collaboration (A. Bazavov *et al.*), *Phys. Lett. B* **795**, 15 (2019), [arXiv:1812.08235 \[hep-lat\]](#).
- [7] A.Y. Kotov, M.P. Lombardo, A. Trunin, *Phys. Lett. B* **823**, 136749 (2021), [arXiv:2105.09842 \[hep-lat\]](#).
- [8] G. Aarts *et al.*, *Phys. Rev. D* **105**, 034504 (2022), [arXiv:2007.04188 \[hep-lat\]](#).
- [9] HotQCD Collaboration (A. Bazavov *et al.*), *Phys. Rev. D* **100**, 094510 (2019), [arXiv:1908.09552 \[hep-lat\]](#).
- [10] M. Dalla Brida *et al.*, *J. High Energy Phys.* **2022**, 034 (2022), [arXiv:2112.05427 \[hep-lat\]](#).
- [11] G. Aarts *et al.*, *Phys. Rev. D* **92**, 014503 (2015), [arXiv:1502.03603 \[hep-lat\]](#).
- [12] G. Aarts *et al.*, *J. High Energy Phys.* **2017**, 034 (2017), [arXiv:1703.09246 \[hep-lat\]](#).
- [13] G. Aarts, C. Allton, D. De Boni, B. Jäger, *Phys. Rev. D* **99**, 074503 (2019), [arXiv:1812.07393 \[hep-lat\]](#).
- [14] M. Marczenko, D. Blaschke, K. Redlich, C. Sasaki, *Phys. Rev. D* **98**, 103021 (2018), [arXiv:1805.06886 \[nucl-th\]](#).
- [15] FASTSUM Collaboration, in preparation, 2022.
- [16] C. Jung *et al.*, *Phys. Rev. D* **95**, 036020 (2017), [arXiv:1610.08754 \[hep-ph\]](#).
- [17] G. Montaña, À. Ramos, L. Tolos, J.M. Torres-Rincon, *Phys. Lett. B* **806**, 135464 (2020), [arXiv:2001.11877 \[hep-ph\]](#); *Phys. Rev. D* **102**, 096020 (2020), [arXiv:2007.12601 \[hep-ph\]](#).