Diagrammatic Monte Carlo

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Sign PROBLEM and Sign PROBLEM

Extensive configurational space

Complexity
$$\propto e^{V/V_0}$$

V is the configuration volume

 V_0 is the correlation volume

$$V/V_0 \sim 10^2$$

Diagrams (intensive configurational space)

Complexity =
$$C(N!)^m$$

 $N \lesssim 10$ is the diagram order

$$m \sim 1$$
, $C \sim 1$

C can be significantly reduced by partial summation

The issue of summability: Can an asymptotic series be regularized?

Dyson's argument: The perturbative series has **zero convergence radius** if changing the sign of interaction renders the system pathological.

BUT

- 1. Fermions can be put on a lattice...
- 2. Bosons can be represented as pairs of fermions...

Feynman's Diagrams

Generic structure of diagrammatic expansions:

$$Q(y) = \sum_{m=0}^{\infty} \sum_{\xi_m} \int D(\xi_m, y, x_1, x_2, \dots, x_m) dx_1 dx_2 \cdots dx_m$$

These functions are visualized with diagrams.

Example:

General principles of Diagrammatic Monte Carlo

Take a diagrammatic series, say the one for polaron Green's function,

and interpret it as a partition function for an ensemble of graphical objects (diagrams). Introduce a Markov process generating ensemble, and calculate corresponding histograms/averages.

The Markov process is organized in the form of pairs of complementary updates. In such a pair, A-B, the update A creates a new graphical element with corresponding continuous variables, while the update B removes the element. For example, A creates a new propagator, while B removes it:



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Balancing diagrammatic Markov process by (generalized) Metropolis-Hastings algorithm

Acceptance ratios for complementary updates A-B

$$R_{A}\left(\vec{X}\right) = \frac{New \, Diagram}{Old \, Diagram} \, \frac{1}{W\left(\vec{X}\right)}$$

$$Arbitrary \, distribution \, function \, for \, generating \, particular \, values \, of \, new \, continuous \, variables \, in \, the \, update \, A$$

$$R_{B}\left(\vec{X}\right) = \frac{New \, Diagram}{Old \, Diagram} \, W\left(\vec{X}\right)$$

Numerical counterpart of analytic bold-line trick: Bold(-line) Diagrammatic Monte Carlo

$$G = G^{(0)} + \cdots + \cdots + \cdots$$

$$G = G^{(0)} + G^{(0)} \Sigma G$$

Dyson equation

Self-consistent determination of G

The Monte Carlo process is <u>asymptotically</u> Markovian

Model of Resonant Fermions

(Good for both ultracold atoms and neutron stars!)

No explicit interactions—just the boundary condition:

$$\forall i, j \text{ at } \left| \mathbf{r}_{\uparrow_i} - \mathbf{r}_{\downarrow_j} \right| \to 0: \qquad \Psi \left(\mathbf{r}_{\uparrow_1}, ..., \mathbf{r}_{\uparrow_N}, \mathbf{r}_{\downarrow_1}, ..., \mathbf{r}_{\downarrow_N} \right) \rightarrow \frac{A}{\left| \mathbf{r}_{\uparrow_i} - \mathbf{r}_{\downarrow_j} \right|} + B, \qquad \frac{B}{A} = c = \text{const}$$

Works whenever $R_0 \ll 1/c$

$$c \gg n^{1/3} \sim k_E \implies \text{BCS regime}$$

$$|c| \ll n^{1/3} \sim k_F \implies \text{unitarity regime}$$

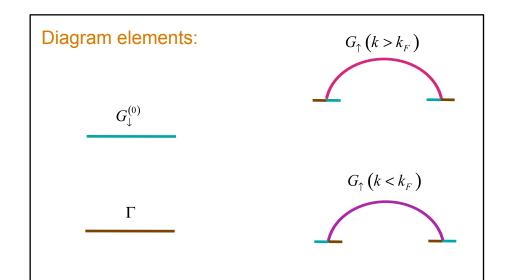
$$-c \gg n^{1/3} \sim k_F \implies \text{BEC regime}$$

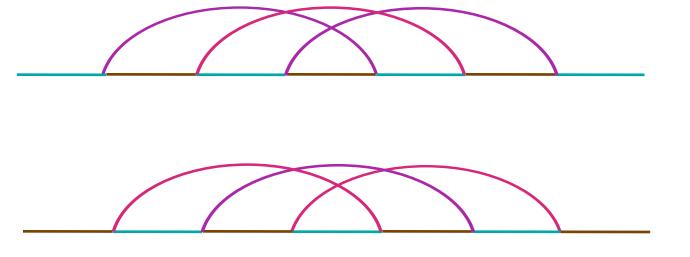
(In two-body problem, the parameter c defines the s-scattering length.)

Resonant Fermipolaron:

Sign alternating divergent series

Prokof'ev and Svistunov, 2007.

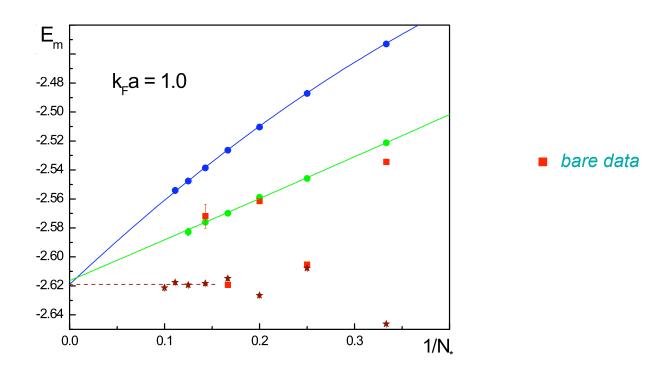




a polaron diagram

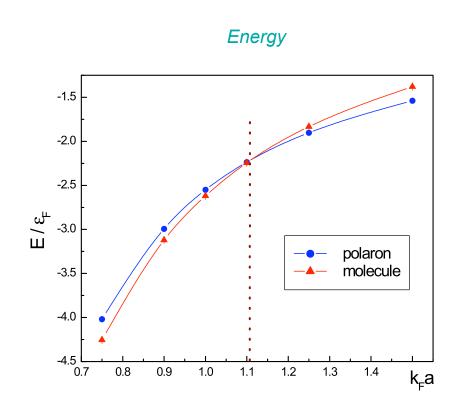
a molecule diagram

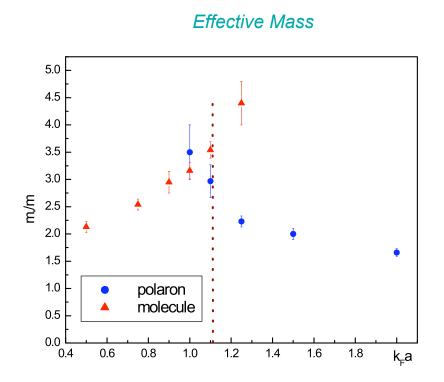
Summability of the divergent series



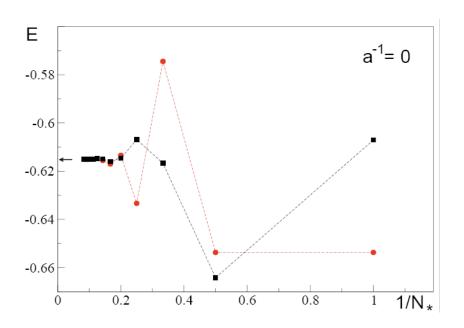
Energy as a function of cutoff parameter for various summation schemes

Answers for fermipolaron





Fermipolaron Energy by Bold Diagrammatic Monte Carlo



- Polaron lines are bold, molecule lines are not.
- Both polaron and molecule lines are bold.

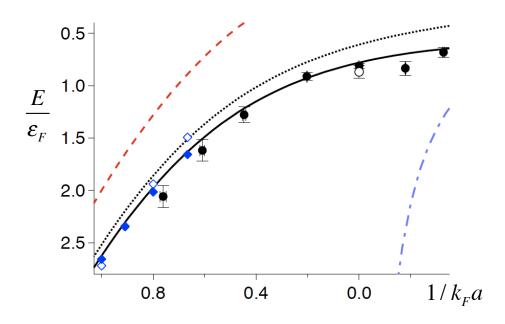
Observation of Fermi Polarons in a Tunable Fermi Liquid of Ultracold Atoms

André Schirotzek, Cheng-Hsun Wu, Ariel Sommer, and Martin W. Zwierlein

Department of Physics, MIT-Harvard Center for Ultracold Atoms, and Research Laboratory of Electronics,

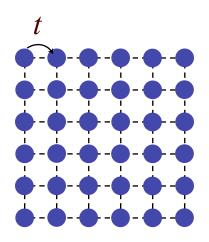
Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

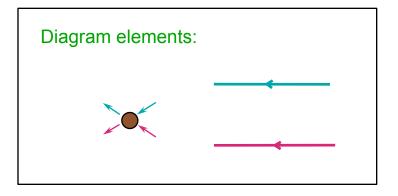
arXiv: 0902.302



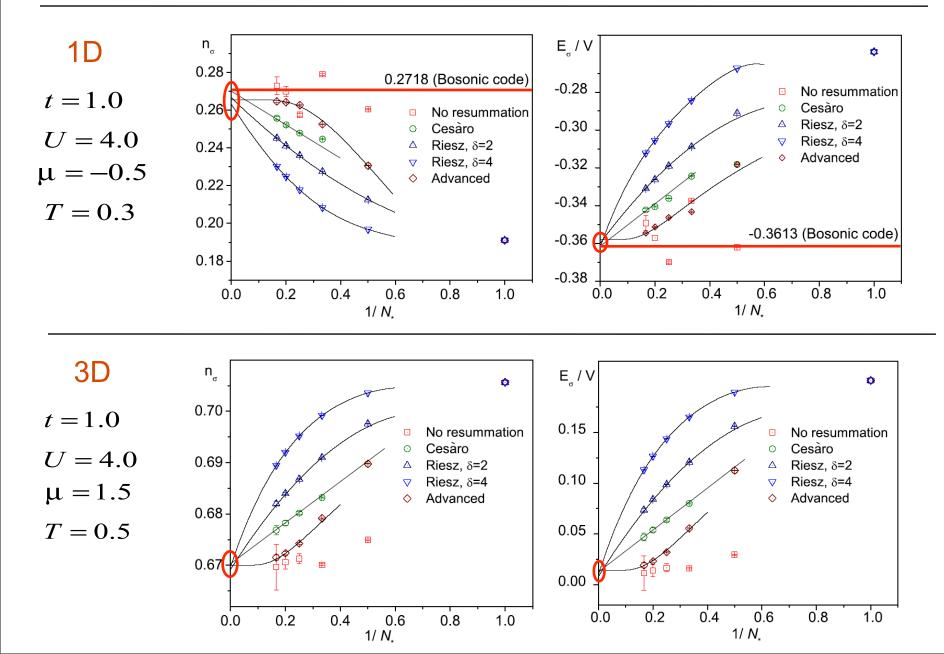
Hubbard model

$$H = -t \sum_{\substack{\langle ij \rangle \\ \sigma = \uparrow, \downarrow}} a_{\sigma i}^{+} a_{\sigma j} + U \sum_{i} n_{\uparrow i} n_{\downarrow i}, \qquad n_{\sigma i} = a_{\sigma i}^{+} a_{\sigma i}$$

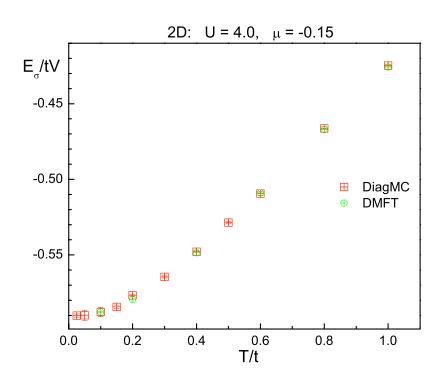




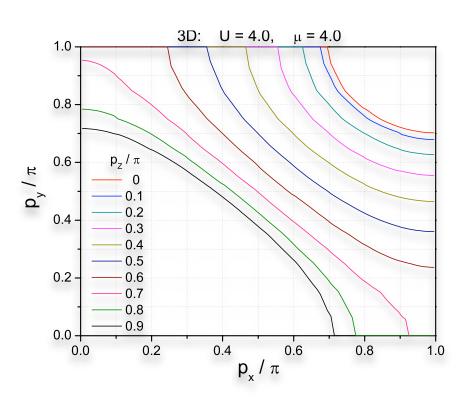
Some Results for the Hubbard Model



Energy per lattice site



Fermi surface



Resonant Fermions at unitarity: Equation of State

